

Heated gas bubbles enrich, crystallize, dry, phosphorylate and encapsulate prebiotic molecules

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Non-equilibrium conditions must have been crucial for the assembly of the first informational polymers of early life, by supporting their formation and continuous enrichment in a long-lasting environment. Here, we explore how gas bubbles in water subjected to a thermal gradient, a likely scenario within crustal mafic rocks on the early Earth, drive a complex, continuous enrichment of prebiotic molecules. RNA precursors, monomers, active ribozymes, oligonucleotides and lipids are shown to (1) cycle between dry and wet states, enabling the central step of RNA phosphorylation, (2) accumulate at the gas-water interface to drastically increase ribozymatic activity, (3) condense into hydrogels, (4) form pure crystals and (5) encapsulate into protecting vesicle aggregates that subsequently undergo fission. These effects occur within less than 30 min. The findings unite, in one location, the physical conditions that were crucial for the chemical emergence of biopolymers. They suggest that heated microbubbles could have hosted the first cycles of molecular evolution.

Life is a non-equilibrium system. Through evolution, life has created a complex protein machinery to maintain the non-equilibrium of crowded molecules inside dividing vesicles. Based on entropy arguments, equilibrium conditions are unlikely to have triggered the evolutionary processes at the origin of life¹. External non-equilibria had to have been provided for the accumulation, encapsulation and replication of the first informational molecules. These can locally reduce entropy, give rise to patterns² and lean the system towards a continuous, dynamic self-organization³. Non-equilibrium dynamics can be found in many fluid systems, including gravity-driven instabilities in the atmosphere⁴, the accumulation of particles in nonlinear flow^{5,6} and shear-dependent platelet activation in blood⁷. Our experiments discuss whether gas–water interfaces in a thermal gradient could have provided such a non-equilibrium setting for the emergence of life on early Earth.

Non-equilibrium systems in the form of heat flows were a very common and simplistic setting found ubiquitously on early Earth⁸. Hydrothermal activity is considered to have been abundant and intimately linked to volcanic activity⁹. Water was thus circulating through the pore space of volcanic rocks, which was formed by magmatic vesiculation (primary origin) and fractures (secondary origin). These systems have been studied as non-equilibrium driving forces for biological molecules in a variety of processes^{10–17}.

Gases originating from degassing of deeper magma bodies percolate through these water-filled pore networks. At shallow depths bubbles are formed by gases dissolved in water and the formation of vapour where sufficient heat is supplied by the hydrothermal system. The bubbles create gas–water interfaces, which

previously have been discussed in connection with atmospheric bubble–aerosol–droplet cycles¹⁸, the adsorption of lipid monolayers and DNA to the interface^{19,20} and the formation of peptide bonds²¹.

In the absence of a temperature gradient, evaporation of a drop of water on a surface exhibits the so-called ‘coffee-ring effect’²². Upon evaporation, molecules in the drop are accumulated at its rim by capillary flow. After complete evaporation, a ring of concentrated material is deposited. In the inverted setting studied here, a gas bubble is immersed in water (Fig. 1) and a temperature gradient drives this process continuously.

Results

Accumulation at the gas–water interface. Experimentally, bubbles were created by filling a 240- μm -thick, corrugated microfluidic chamber with solution (Supplementary Fig. 1). As the solution could not fill all the cavities, pinned gas bubbles were created. At higher temperatures, bubbles were found to also form spontaneously anywhere in the system as a result of outgassing. These bubbles were often not restricted by their surrounding geometry and moved along the heated surface.

Accumulation of molecules at these heated gas–water interfaces is caused by a continuous evaporation–recondensation water cycle. The interfaces are held in a constant non-equilibrium state, leading to a steady-state coffee-ring effect that does not end in a fully dry state. Here, we observe six physico-chemical processes, which have all individually been suggested to be relevant for the emergence of prebiotic evolution and are co-located in a single non-equilibrium system: (1) enhanced catalytic activity of ribozymes, amplified by the accumulation of oligonucleotides and ions, (2) condensation of

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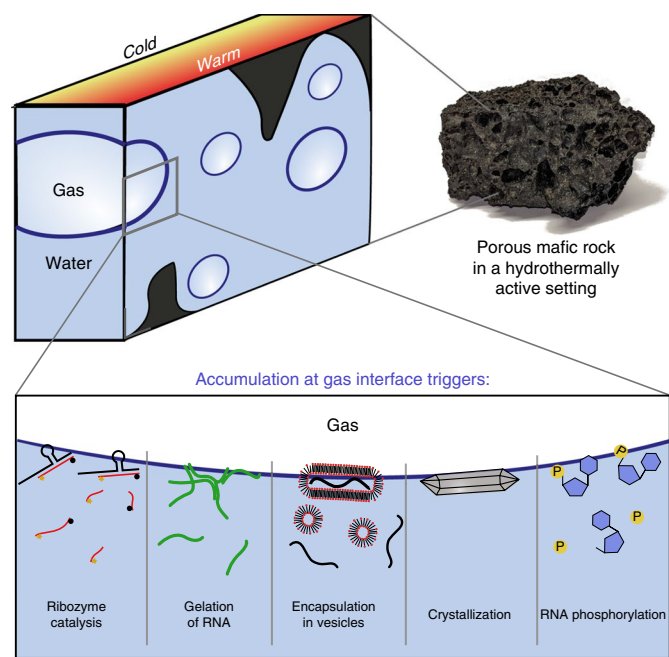


Fig. 1 | DNA accumulation at gas bubbles in a thermal gradient. Volcanic rocks in shallow hydrothermal settings are subjected to water cycles in their pore spaces, which are of both primary (magmatic vesicles) and secondary (fractures) origin. Gases, originating from magma degassing at depth, percolate through the water. Heat supplied by the hydrothermal system causes vaporization. At the gas microbubbles, molecules are accumulated by the continuous capillary flow on the warmer side of the gas–water interface. As shown experimentally, this environment can enhance the catalytic activity of ribozymes, trigger the formation of a hydrogel from self-complementary RNA, encapsulate oligonucleotides such as aptamers in vesicle aggregates, trigger their subsequent fission, drive the crystallization of ribose aminooxazoline (RAO)—a prebiotic RNA precursor—and initiate the phosphorylation of RNA nucleosides.

self-complementary RNA 36mers into millimetre-sized hydrogels, (3) vesicle aggregation at the bubble interface along with encapsulation of oligonucleotides in aqueous phases with up to 18-fold enhanced concentration, (4) fission of the vesicle structures in the adjacent micro-convection, (5) formation of euhedral 300 μm crystals from the RNA precursor ribose aminooxazoline (RAO) around bubbles (which also act as seeds for new bubbles) and (6) dry–wet cycles enhancing, for example, the phosphorylation of nucleosides, created by fluctuating and moving interfaces (Fig. 1 and Supplementary Videos 1–6). All six mechanisms were established within 30 min and, importantly, operated in continuous contact with bulk water.

Here, length- and temperature-dependent accumulations were measured at low-salt conditions (0.1-fold PBS buffer: 13.7 mM NaCl, 0.27 mM KCl, 1 mM phosphate buffer). However, DNA and RNA gelation measurements were performed equally well under physiological conditions (1-fold PBS buffer: 137 mM NaCl, 2.7 mM KCl, 10 mM phosphate buffer).

To observe the dynamics at the interface, bubbles were created as described above, using fluorescently labelled molecules in solution (Fig. 2a(i)). The front and back sides of the chamber were heated and cooled, respectively, to generate a temperature gradient. The system was monitored through the warm side using a fluorescence microscope. The optical axis ran along the temperature gradient (Supplementary Figs. 1 and 2) and the quantitative fluorescence was captured with a charge-coupled device (CCD) camera.

Initially, the chamber was filled with a solution of 200 nM 6-Carboxyfluorescein (FAM)-labelled 132-base single-stranded DNA (ssDNA) oligomer in 0.1-fold PBS buffer. When no temperature gradient was applied to the system ($T_{\text{warm}} = T_{\text{cold}} = 10^\circ\text{C}$), we observed no accumulation of DNA near the gas–water interface. The fluorescence signal exhibited a constant small peak at the observed interface, possibly due to a slight adsorption of the DNA to the gas–water interface (Fig. 2a(i), 0 s). Heating one side of the chamber ($T_{\text{warm}} = 30^\circ\text{C}$, $T_{\text{cold}} = 10^\circ\text{C}$) resulted in the rapid accumulation of DNA in a small area on the warm side at the contact line (Fig. 2a(ii), dashed red box, Supplementary Video 1). The chamber-averaged fluorescence at the contact line increased within 6 min by ~ 12 -fold compared to the bulk fluorescence (Fig. 2b). No accumulation was observed on the cold side.

We calculated the local concentration at the contact line from the ratio of the meniscus and bulk fluorescence and the geometry of the curved gas–water interface. Because the fluorescence was averaged over the chamber by the microscope objectives, a 60-fold higher concentration in addition to the higher fluorescence is inferred due to the thinner size of the accumulation region ($\sim 4 \mu\text{m}$) compared to the 240- μm -wide chamber. Therefore, from the observed 12-fold increase in fluorescence, we estimated a concentration increase by a factor of 700, corresponding to 140 μM DNA concentration in the meniscus when starting from a 200 nM bulk solution (Fig. 2c). Simulations suggested that without thermophoresis this accumulation would be only slightly higher, showing that it does not play a significant role in the accumulation process (Fig. 2c, red dotted line).

Further analysis from experiment and theory showed that the accumulation was caused by the focused evaporation of water at the tip of the meniscus²³ (Fig. 2d, orange). A continuous flow of water into the meniscus dragged the molecules with it; because they could not evaporate, they could only escape by diffusion against the one-way capillary flow. The flow was visualized by filling the chamber with a suspension of 200-nm-diameter FAM-labelled polystyrene beads. By particle tracking, we measured the velocity profile in the meniscus (Fig. 2d,e and Supplementary Video 1). Beads moved towards the accumulation region near the hot side of the chamber at the contact line.

We attribute this to capillary flow, which superseded the comparatively weak bulk buoyant convection (Fig. 2e). Temperature gradients have also been demonstrated to create Marangoni flows²³, in which water is drawn from the warm to the cold side of an interface due to a surface tension gradient. This was difficult to observe because the main temperature gradient was along the viewing axis. However, we observed strong lateral flows at the interface when accumulating vesicles. We attribute these to lateral Marangoni flows and therefore assume also a combination of Marangoni flows and convection along the interface (Fig. 2d, green). Without this flow, our simulation predicts a larger accumulation (Fig. 2c, blue dotted line).

The water that evaporated near the warm chamber wall was found to condense on the cold wall, forming small water droplets. As a result, the gas bubble had less space, expanded, and moved the gas–water interface. Once the condensed ‘rain’ droplets had grown, merged and re-entered into the bulk solution, driven by surface tension, the interface moved back to its initial position (Supplementary Video 1). These fluctuations of the contact line triggered a drying and recondensation of the molecules at the location where the capillary flow initially accumulated them.

A fluid dynamics model was used to describe the main features of the accumulation. We considered a two-dimensional (2D) section perpendicular to the interface. Because the width of the channel was smaller than the capillary length for the water and gas²⁴, the interface geometry was approximated by an arc of a circle.

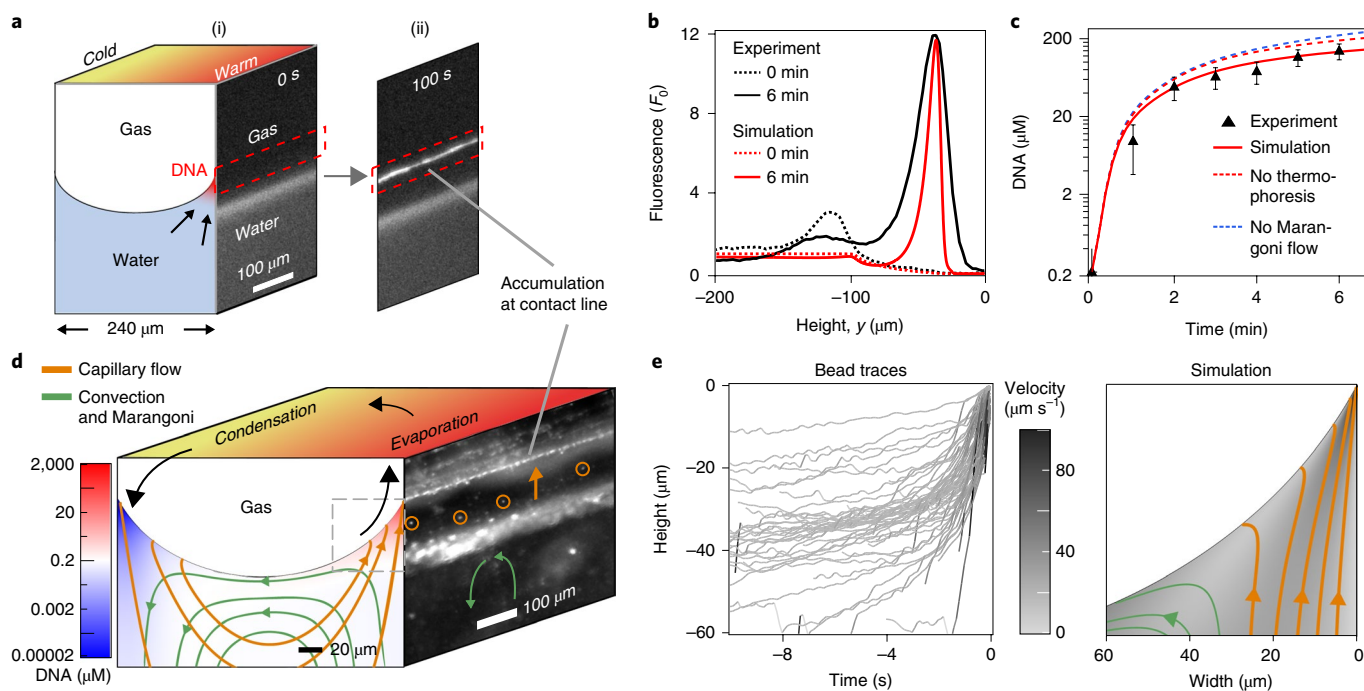


Fig. 2 | DNA accumulation by capillary flow. **a**, (i) Gas–water interface imaged by fluorescence microscopy with the contact line area outlined with a red dashed line. The initial higher fluorescence intensity at the contact line originated from adsorbed DNA at the gas–water interface. (ii) After 100 s, the 132mer DNA has accumulated in water near the contact line at the warm temperature side of the chamber ($T_{\text{warm}} = 30\text{ }^{\circ}\text{C}$, $T_{\text{cold}} = 10\text{ }^{\circ}\text{C}$, Supplementary Video 1). **b**, The fluorescence profile reached a 12-fold increase as compared to the bulk fluorescence within 6 min. **c**, Due to the confined meniscus geometry, the accumulated concentration was significantly higher. The 12-fold increase in fluorescence corresponded to a 700-fold increase in DNA concentration, consistent with simulation results (red solid line). Without thermophoresis (red dotted line) or Marangoni flows (blue dotted line), simulations predict a slightly higher accumulation. Error bars were estimated from fluorescence analysis averaged over an $\sim 200\text{ }\mu\text{m}$ interface width at three positions along the contact line from the shown example measurement. **d**, Fluid flow near the contact line measured by single-particle tracking (Supplementary Video 1). Capillary flow (orange) pulled the beads upwards toward the contact line as water mainly evaporated at the tip of the meniscus. Marangoni and convection flows (green) shuttled the beads between the hot and cold sides. Superposed is the logarithmic concentration profile obtained from the simulation for the accumulation in **a**. **e**, Single-particle tracking of the capillary flow. Peak flow velocities reached $50\text{ }\mu\text{m s}^{-1}$ in the last second of the water flow before its evaporation (Supplementary Fig. 3). $T_{\text{cold}} = 10\text{ }^{\circ}\text{C}$, $T_{\text{warm}} = 40\text{ }^{\circ}\text{C}$.

The contact line of the interface was pinned in the simulation^{25,26}, motivated by the observation that the accumulation kinetics was generally faster than the fluctuating movement of the interface.

The model superposed four water flows, providing the boundary condition for the accumulation of DNA: (1) capillary flow at the meniscus, (2) diffusion of water vapour between the interface and the gas bubble, (3) convection of water and (4) Marangoni flow along the interface. The relative strength of the Marangoni flow—a free parameter due to the unknown presence of surface-active molecules²⁷—was adjusted to fit the velocities measured in the experiment.

The interplay of all four flows led to the accumulation of molecules at the meniscus. As only water evaporated on the warm side, dissolved molecules were continuously dragged towards the contact line, where their concentration depended on back-diffusion and the speed of the capillary flow. Convection and Marangoni flow provided a constant cycling of water and new material towards the accumulation region.

DNA accumulation was measured experimentally for various temperature differences and DNA lengths (Supplementary Fig. 4). The accumulation rose with increasing temperature difference (ΔT), reaching a 4,000-fold increase for $\Delta T = 40\text{ }^{\circ}\text{C}$. We also found that smaller (15mer) DNA molecules accumulated three times less effectively than larger (132mer) DNA, which was attributed to their higher diffusion coefficient. The model predicted a multi-fold

accumulation of mono- and divalent ions (Supplementary Fig. 4), resulting in a higher salt concentration at the meniscus.

Enhanced ribozyme catalysis at the interface. The accumulation of larger biomolecules as well as ions at the interface makes it a powerful mechanism to enhance the activity of functional nucleic acids. To test this, we monitored the activity of the Hammerhead ribozyme^{28,29}, which cleaves a 12mer RNA substrate strand. The substrate and magnesium concentration determined its activity³⁰, and both were accumulated in a $30\text{ }^{\circ}\text{C}$ gradient ($T_{\text{warm}} = 10\text{ }^{\circ}\text{C}$, $T_{\text{cold}} = 40\text{ }^{\circ}\text{C}$) at low bulk concentrations ($0.1\text{ }\mu\text{M}$ Hammerhead, $0.5\text{ }\mu\text{M}$ substrate, 0.4 mM MgCl_2 , Fig. 3a). As a control, chambers without gas interfaces were studied. A FAM dye and a black hole quencher were attached on opposite sides of the substrate, inhibiting the fluorescence of the FAM dye. On cleavage, the dye was not quenched anymore and could be detected by fluorescence microscopy. Figure 3a,b and Supplementary Video 2 show the average fluorescence of the chamber over time. The substrate strands were cleaved predominantly at the interface, as seen by the rise of fluorescence there. From here, the cleaved strands were frequently ejected into the bulk solution. After 25 min, samples were extracted from the bulk fluid and analysed by polyacrylamide gel electrophoresis (PAGE, Fig. 3c). We detected up to 50% concentration of the cleaved substrate (bottom band, Supplementary Fig. 5) in a chamber with bubbles.

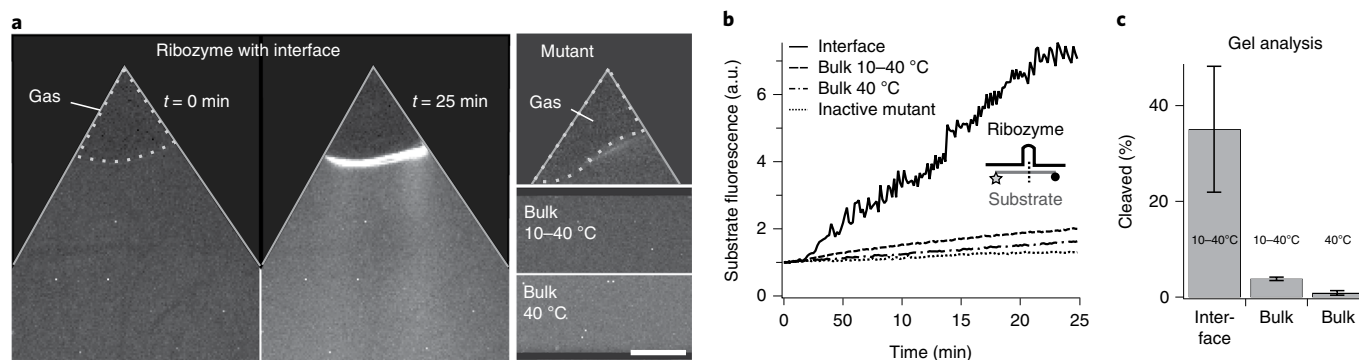


Fig. 3 | Ribozyme catalysis triggered by an interface. Fluorescence microscopy and PAGE analysis of Hammerhead ribozyme activity. **a**, Hammerhead in a chamber with a gas–water interface after 0 and 25 min (left and right, respectively; $T_{\text{warm}} = 40^\circ\text{C}$, $T_{\text{cold}} = 10^\circ\text{C}$, $0.1\ \mu\text{M}$ ribozyme, $0.5\ \mu\text{M}$ substrate, $0.4\ \text{mM}$ MgCl_2). Fluorescence increased strongly at the interface over time (Supplementary Video 2). Without the interface (bottom right), the Hammerhead shows little activity, both in the same temperature difference and at a homogeneous 40°C ; an inactive mutant (top right) showed no significant increase in fluorescence. Scale bar, $500\ \mu\text{m}$. **b**, Overall fluorescence in the bulk fluid from **a** over time. The illustration shows the ribozyme (black) and substrate (grey) with the cleavage site (dotted line), dye (star) and quencher (black circle). **c**, PAGE analysis results with and without the interface after 25 min (gels in Supplementary Fig. 5). We found a significant increase in cleaved product with the interface ($35 \pm 13\%$, five measurements) compared to bulk samples in the same temperature gradient ($3.79 \pm 0.35\%$, three measurements) or at 40°C ($0.86 \pm 0.49\%$, three measurements). Error bars show s.d. from normalized gel band intensities.

In the same temperature gradient, but without interfaces, only 3.8% of the substrate was cleaved. This decreased even further when the bulk temperature was set homogeneously to 40°C , where only very little activity could be observed. When using an inactive mutant version of the ribozyme, the fluorescence did not increase noticeably (Fig. 3a, top right; Fig. 3b, dotted line), demonstrating that it did not originate from enhanced hydrolysis at the interface. This shows that the accumulation mechanism can enhance the catalytic activity of ribozymes, while maintaining low-salt conditions in the bulk solution. It increases the turnover—the number of strands cleaved per ribozyme—by the accumulation of substrates and ions.

RNA/DNA gelation. The length selectivity observed in nucleic acid accumulation could increase the concentration of self-complementary oligonucleotides to the point that hydrogel-forming concentrations could be reached (Fig. 4a). Self-complementary strands were found to form a macroscopic, millimetre-sized hydrogel, a process previously shown for DNA in a thermophoretic accumulation chamber¹⁵. We observed the formation of hydrogels for self-complementary DNA and RNA and both for GC-only (Fig. 4 and Supplementary Video 3) and AT-only sequences (Supplementary Fig. 6). All oligonucleotides were end-labelled and HPLC purified to minimize the presence of free dye in the experiments.

Starting from uniformly distributed DNA in bulk solution (Fig. 4a(i), initial concentration $10\ \mu\text{M}$; 1-fold PBS buffer), the DNA was quickly accumulated at the interface once the temperature gradient was established. Within 8 min, a hydrogel had formed, after which it quickly detached from the interface and entrained in the convection flow (Fig. 4a(ii)). The hydrogel nature of the DNA was checked by increasing the temperature from 30°C , below the melting point of the self-complementary sequences, to 70°C , where the hydrogel dissolved into the convection flow (Fig. 4a(iii)).

The formation of hydrogels from self-complementary RNA was also observed. We co-accumulated two different sequences in a 20°C temperature difference (Fig. 4b). The red fluorescence channel monitored a 36mer non-complementary ssRNA strand and the green fluorescence channel a self-complementary GC-only 36mer ssRNA strand with three self-complementary binding sites. Based on simulations (Nupack, www.nupack.org), the hydrogel-forming strands bind to each other and form a network of polymers (Fig. 4b,

right). Over the course of the experiment, both the non-complementary and self-complementary strands accumulated at the same interface. However, after 21 min, we only observed a hydrogel for the self-complementary green RNA. No hydrogel could be observed for the non-complementary red RNA, which also accumulated near the surface but was not forming large-scale structures (Supplementary Video 3). Replacing the GC-only strand with a 60mer AU-only RNA with similar self-complementarity gave the same results (Supplementary Fig. 6). This demonstrates that gel formation and separation of the strands is not dominated by G-quadruplex formation.

The red and green strands separated macroscopically based only on their sequence. The self-complementary strands remained at a local high concentration in the hydrogel, which offered reduced hydrolysis rates due to its predominantly double-stranded nature. This sequence-selective gelation was also similarly found for DNA (Supplementary Fig. 6).

DNA encapsulation in vesicle aggregates. A key requirement for the emergence of cellular life is the encapsulation of molecules at increased concentration relative to their more dilute external environment. Fatty acids are potential candidates that could separate nucleic acids in vesicles, possibly incorporating phospholipids into their membranes over time^{31,32}. For the encapsulation, three autonomous processes need to occur: (1) accumulation of oligonucleotides to meaningful concentrations, (2) accumulation of vesicles to trigger their aggregation or fusion and (3) the combination of both in one location to encapsulate oligonucleotides into vesicular structures. Here, we show that heated gas–water interfaces could fulfil these requirements. The accumulated vesicles do not necessarily form larger, round vesicles, but aggregate. However, we demonstrate in the following that these aggregates enclose DNA and RNA in an aqueous phase, allowing their binding and folding. Shear flows close to the interface as well as in the convection flow were shown to divide these aggregates to form smaller and more round structures.

To introduce lipids in a homogeneous manner, the chamber was filled with small 100-nm-sized vesicles prepared from a 10mM oleic acid solution ($0.2\ \text{M}$ Na-bicine, $1\ \text{mM}$ EDTA, pH 8.5) and $2\ \mu\text{M}$ DNA. These initially small vesicles appeared as a continuous

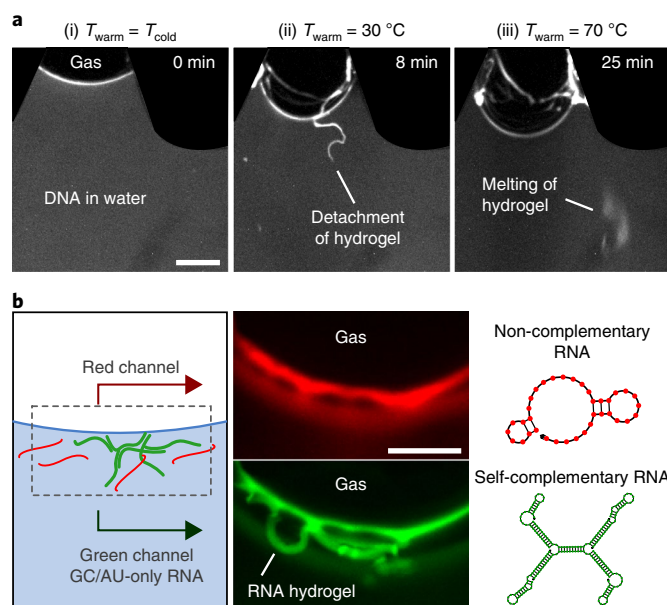


Fig. 4 | Sequence-selective gelation of RNA and DNA. Fluorescence microscopy of DNA and RNA revealed the formation of hydrogels at the interface.

a, Gelation of DNA. (i) Before applying the temperature gradient, DNA accumulated at the interface only due to slight surface adsorption to the gas–water interface. (ii) After applying the temperature gradient, a DNA hydrogel formed and detached from the interface (Supplementary Video 3). (iii) At $T_{\text{warm}} = 70\text{ }^{\circ}\text{C}$, the hydrogel melted and redistributed the DNA back into the bulk fluid. We estimated the DNA concentration in the hydrogel to be $100\text{ }\mu\text{M}$. Scale bar, $500\text{ }\mu\text{m}$. **b**, Gelation of RNA. In a single experiment, non-complementary 36mer ssRNA (red) was accumulated with self-complementary GC-only 36mer ssRNA (green). Both strands accumulated at the interface ($T_{\text{warm}} = 30\text{ }^{\circ}\text{C}$, $T_{\text{cold}} = 10\text{ }^{\circ}\text{C}$), but only the self-complementary RNA formed an elongated, fibrous hydrogel. The same behaviour was found for DNA and AU-only RNA (Supplementary Fig. 6). Scale bar, $125\text{ }\mu\text{m}$.

background (Fig. 5a,b and Supplementary Video 4). After turning on the temperature gradient, we observed the accumulation of these vesicles together with DNA within 10 min at the interface ($T_{\text{warm}} = 70\text{ }^{\circ}\text{C}$, $T_{\text{cold}} = 10\text{ }^{\circ}\text{C}$). The vesicles aggregated together and formed larger clusters. It should be noted that this aggregation and cluster formation were strongly increased if $\sim 0.1\%$ 1,2-dioleoyl-*sn*-glycero-3-phosphoethanolamine was added to the lipids (present here to also label the lipids). Interestingly, the co-accumulated DNA was encapsulated into these vesicle aggregates, in which we found an up to 18-fold increase in oligonucleotide concentration compared to the bulk solution. These aggregates were shuttled into the convection and frequently formed thread-like structures. Close to the interface, we observed strong flows (Supplementary Video 4), which we attribute to Marangoni flows. These could originate from lateral temperature gradients across the chamber due to inhomogeneous heating or differences in the thermal conductivity of the chamber material and water. Aggregates were observed to divide and split into smaller compartments in these flows and convection (Supplementary Fig. 4).

Vesicles formed from the phospholipid 1,2-dioleoyl-*sn*-glycero-3-phosphocholine (DOPC) (3.6 mM DOPC, 0.2 M Na-bicine, 1 mM EDTA, $\text{pH} 8.5$) showed a similar behaviour, but they aggregated more strongly at the interface, and externally applied flow across the microfluidics was sometimes necessary to remove them from the interface. The accumulated 100-nm -sized vesicles formed larger and more spherical structures compared to the oleic acid aggregates, and encapsulated DNA equally well. They also underwent fission due to shear stress in the convection flow (Fig. 5c and Supplementary Video 4).

The vesicle aggregates (oleic acid or DOPC) concentrated around smaller gas bubbles ($150\text{ }\mu\text{m}$ in diameter), inducing a clustering of DNA (Supplementary Fig. 7). This was not observed in the absence of lipids. Lipid vesicles therefore significantly enriched the local DNA concentration.

The co-location of DNA and lipids raises the question of whether the oligonucleotides are in an aqueous phase that allows, for example, the folding of RNA or binding of DNA, and whether these compartments are protected from their surrounding. To test the former, we accumulated the RNA aptamer ‘Broccoli’³³ (Fig. 5d), which folds around the fluorophore DFHBI-1T and increases its fluorescence, with DOPC vesicles ($1\text{ }\mu\text{M}$ aptamer, $10\text{ }\mu\text{M}$ DFHBI-1T, 3.6 mM DOPC, 50 mM HEPES, $\text{pH} 7.6$, 100 mM KCl, 1 mM MgCl_2). Here, we used $T_{\text{warm}} = 40\text{ }^{\circ}\text{C}$ and $T_{\text{cold}} = 10\text{ }^{\circ}\text{C}$ to avoid RNA and fluorophore degradation. Again, the RNA, fluorophore and vesicles accumulated at the interface, leading to the formation of aggregates that were visible both in the lipid as well as the fluorophore colour channel. Replacing the Broccoli aptamer with a non-binding RNA strand led to a more than 100-fold reduced fluorescence of the fluorophore. This shows that the aptamer was folded inside the aggregates.

To demonstrate that the accumulated material was protected inside the aggregates, a 72mer double-stranded DNA (dsDNA) was accumulated with DOPC vesicles (Fig. 5e, $7.1\text{ }\mu\text{M}$ DNA, 3.6 mM DOPC, 0.2 M Na-glycineamide $\text{pH} 8.5$, 6 mM MgCl_2 , 1 mM CaCl_2). Each pair of dsDNA strands thereby contained a FAM dye on one strand and a 5-Carboxy-X Rhodamin (ROX) dye opposite to it on the other strand. These dyes form a Förster resonance energy transfer (FRET) pair, in which the excited FAM dye can transfer energy to the ROX dye, which then fluoresces. As this energy transfer works only in close proximity of the dyes, we could use it to measure the amount of dsDNA inside the vesicle aggregates. The higher the signal—between 0 and 1—the more dsDNA was present. After formation of aggregates at the interface, most of the solution was extracted and DNase I was added (0.5 units per $\sim 10\text{ }\mu\text{l}$). The solution was then put back into the chamber and the FRET signal was observed at $37\text{ }^{\circ}\text{C}$. The DNase digested the DNA strands outside the aggregates, where the signal quickly dropped (Fig. 5e, Supplementary Video 4 and Supplementary Fig. 7). The FRET signal of the aggregates reduced only slightly, probably as a consequence of the digestion

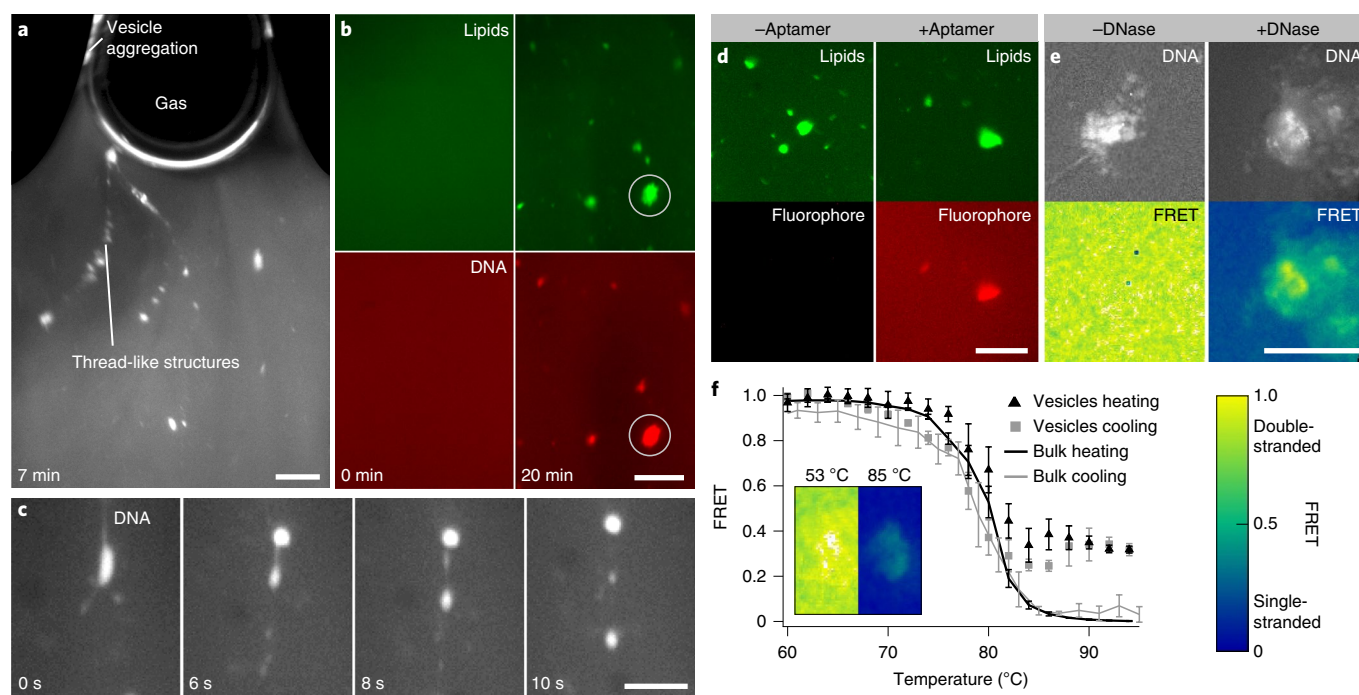


Fig. 5 | DNA and RNA encapsulation and protection in vesicles formed at the interface. Fluorescence microscopy images of DNA/RNA and lipid colour channels. **a**, Lipid channel, with 100 nm sized oleic acid vesicles accumulated at the gas-water interface. These formed aggregates of vesicles that were ejected into the bulk and formed thread-like structures (Supplementary Video 4). **b**, Comparison of lipid and DNA colour channels (green, lipid; red, DNA), showing that lipids and DNA co-localize, meaning that the 100 nm-sized vesicles have formed larger clusters and encapsulated DNA at up to 18-fold enhanced concentration compared to the bulk. **c**, DNA channel, with DOPC vesicles accumulated under the same conditions and also encapsulated DNA at enhanced concentrations. Here, a fission of the DOPC vesicles containing DNA was observed in the shear flow. **d**, Encapsulation of fluorophore DFHBI-1T with a non-binding RNA (left) and the Broccoli aptamer (right), which increases its fluorescence. It shows that RNA has folded in an active conformation inside the vesicles. **e**, FRET analysis of dsDNA in aggregates before (left) and after (right) addition of DNase I. DNA inside the aggregates was protected from the DNase (high FRET), whereas it was digested in the bulk solution (low FRET, Supplementary Video 4). **f**, FRET melting curve of dsDNA inside and outside the aggregates, starting from 1.0 (all DNA double-stranded) at 60 °C to 0 (all DNA single-stranded) at 95 °C outside and 0.3 inside the aggregates. Error bars indicate s.d. measured for a larger area of normalized signal of the shown example measurement. Scale bars, 200 μm (**a,b**), 100 μm (**c**), 250 μm (**d,e**).

of solution above/below or DNA sticking to their surface. We thus show that the aggregates protect the molecules inside them.

The FRET signal was also used to observe the melting of dsDNA inside the chamber (Fig. 5f). After accumulation (7.1 μM DNA, 3.6 mM DOPC, 0.2 M Na-glycineamide pH 8.5, 11 mM NaCl, 0.22 mM KCl, 0.8 mM phosphate buffer), the chamber was heated to 95 °C, during which the FRET signal reduced to 0 in the bulk and ~ 0.3 in the aggregates. The remaining signal could stem from aggregated DNA or DNA that is enclosed and stabilized by lipids in a way that did not allow the strands to fully unbind. After cooling, the FRET signal returned to 1 for both aggregates and the bulk solution. The combination of aptamer and FRET analysis demonstrates that oligonucleotides are encapsulated inside the aggregates in an aqueous phase that allows them to melt and fold.

The above experiments assumed the presence of uniformly 100-nm-sized vesicles at the beginning. We also explored the behaviour of the system when it initially contained a range of vesicle sizes, up to $\sim 30 \mu\text{m}$ in diameter (Supplementary Fig. 8). These exhibited a similar, but often slower accumulation behaviour. Within 20 min, the system again started to form vesicle clusters that contained enhanced DNA concentrations. The formation of oleic acid aggregates was also observed in the absence of DNA, indicating that they were not the result of a DNA/lipid interaction (Supplementary Fig. 8).

Crystallization at gas bubbles. The building blocks for the synthesis of single nucleotides, such as the prebiotic RNA precursor ribose

aminoxazoline (RAO), accumulated near gas bubbles to concentrations that triggered its crystallization. For RAO, a crystallization is of fundamental interest because it can be both diastereoisomerically purified by selective sequential crystallization of its precursors and enantiomerically enriched by conglomerate crystallization, where the two enantiomers (D- and L-) of RAO crystallize into discrete independent domains^{34,35}. To trigger controlled crystal growth, RAO would need to be accumulated slowly around a growing bubble. In previous experiments, bubbles were artificially created. Here we used elevated temperatures (70 °C warm side, 10 °C cold side) to trigger the spontaneous formation of a bubble. Subsequently, RAO accumulated and crystallized around it (Fig. 6a,b). Therefore, a 40 mM solution of D-RAO, 2.8-fold below the saturation concentration at $T_{\text{warm}} = 70 \text{ }^\circ\text{C}$ ($\sim 110 \text{ mM}$), was used to fill a chamber without a corrugated geometry and did not create a gas-water interface. For the crystallization to occur at the warm side of the chamber, RAO would have needed to accumulate several fold to overcome the nucleation energy barrier³⁶. We monitored the fluorescence of 1 μM Cy5 added to the solution, which co-accumulated at the gas-water interface but was not incorporated into the crystals.

An initially small bubble formed on the warm side and accumulated RAO around it. Within 40 min, the bubble grew while continuously increasing the RAO concentration at its warm side. The crystal shown in Fig. 6a was found at the location where the bubble had formed. No crystals were found on the cold side of the chamber. An X-ray crystal structure determination confirmed that crystals grown on the warm side were indeed D-RAO (Fig. 6c).

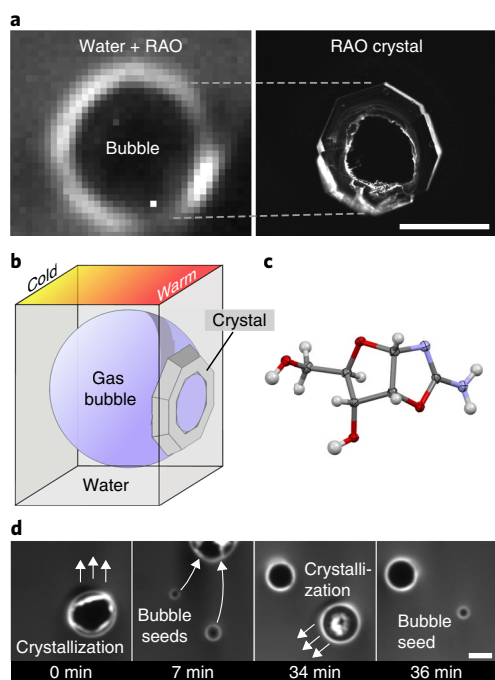


Fig. 6 | Crystallization and bubble movement. **a**, RAO accumulated from 40 mM bulk concentration around a growing bubble at the warm side of the chamber (70 °C). After 37 min, an RAO crystal formed below the bubble at the water–gas–solid interface, as indicated. **b**, The crystal shape suggests that it formed successively as the bubble grew, consistent with simulation results (Supplementary Fig. 9). **c**, X-ray crystal structure analysis of a D-RAO crystal grown at a gas–water interface. **d**, Moving bubbles leave behind crystals (7 min) that act as seeds for new bubbles (34 min, Supplementary Video 4). Scale bars, 250 μm .

Interestingly, the formation of RAO crystals led to the later re-formation of gas bubbles near small cavities—a known and well described process of heterogeneous nucleation in supersaturated solutions³⁶. We observed this crystal-induced bubble formation in experiments at high $T_{\text{warm}} = 70\text{ }^{\circ}\text{C}$ (Fig. 6d and Supplementary Video 4). When a gas bubble (around which RAO crystals had formed) moved away, it left behind accumulated material, possibly crystal cavities still filled with gas, which facilitated the growth of new daughter bubbles at the same location. When these daughter bubbles increased in size sufficiently, they started to accumulate RAO, formed crystals and eventually moved away. The remaining crystals again hosted the growth of new bubbles and began the bubble-induced crystallization cycle again.

Dry–wet cycles and phosphorylation by moving interfaces.

If bubbles were not confined by their geometry, they tended to move upwards in the chamber (Fig. 7a) due to buoyancy forces. This movement led to a continuous cycling of dry–wet conditions (Fig. 7a), as accumulated material close to the interface entered the bubble, dried and was rehydrated when the bubble moved away. At the same time, accumulated material at the trailing edge was dragged along if the bubble moved slowly enough, keeping high molecule concentrations in the vicinity of the interface (Supplementary Video 6).

At a hydrophilic surface material such as silicon dioxide (quartz), used so far on the cold side in the experiments, the contact angles of the condensed droplets were small, and they re-entered quickly into the bulk water, leading to only little movement of the gas–water interface. To trigger many wet–dry cycles, we placed a Teflon foil

on the warm and cold sides of the chamber. This led to the formation of many droplets on the cold side (Fig. 7b, small round structures inside the dark gas region, Supplementary Video 6), which, if large enough, could also be in contact with the warm side. In this setting, we tested the phosphorylation of nucleosides, a reaction that usually requires dry conditions at elevated temperatures. The reaction is known to be most effective when the solution is dried at 100 °C, a scenario hard to reconcile with typical RNA-world conditions^{37,38}. We observed a 20 times more efficient phosphorylation of cytidine nucleosides (Fig. 7c, $T_{\text{warm}} = 60\text{ }^{\circ}\text{C}$, $T_{\text{cold}} = 37\text{ }^{\circ}\text{C}$, 240 mM cytidine and ammonium dihydrogen phosphate, 2.4 M urea, similar to the phosphorylation used in ref.³⁷) compared to bulk water at average temperatures of 50 °C and 60 °C. The found enhanced in situ phosphorylation reaction would improve the recycling of hydrolysed RNA.

Discussion

We have found a general accumulation mechanism of molecules and small vesicles at gas bubbles subjected to heat flow in water. A temperature gradient across a gas–water interface created a continuous process of evaporation and condensation of water between the warm and cold sides. This moved molecules towards the bubble interface and increased their concentration by several orders of magnitude, depending on their diffusion coefficients. Because the contact line between gas and water was mobile, dry–wet cycles were created, in many conditions as often as twice per minute. The molecules studied here have been discussed as prebiotic candidates before the emergence of life^{35,37,39}. Fluorescently labelled analogues were used to probe their concentration. We did not elaborate on the types of gas used in this study, because, due to their low concentrations, they did not modify the surface tension and water evaporation significantly⁴⁰.

Our simulation captured the basic characteristics of this accumulation and validated the experimental results. A 4,000-fold DNA accumulation was reached both in experiment and simulation. In comparison with thermophoretic traps, with which the setting could be combined, the accumulation at the interface occurred significantly faster, on the timescale of minutes rather than hours¹¹.

We used different temperature gradients adapted to the different scenarios. The magnitude of accumulation depended strongly on the applied temperature difference (Supplementary Fig. 4), which was kept at high values in Figs. 3 and 4 (10–40 °C and 10–30 °C, respectively) to accumulate RNA strongly, but still under cold conditions to keep hydrolysis insignificant. Vesicles aggregated under similar conditions (Fig. 5d), but with faster convective flows from a higher temperature difference (10–70 °C), shuttled aggregated vesicles more efficiently away from the interface for downstream analysis. The creation of crystals from RAO (Fig. 6) required the larger temperature gradients required for an efficient accumulation of the comparably small molecules. The phosphorylation chemistry of monomers (Fig. 7) profited from an overall enhanced temperature.

The co-accumulation of small ions offers additional reactivity. Divalent salts such as Mg^{2+} were predicted by our model to accumulate by a factor of 4–5 for temperature differences of 20–30 °C. In high-salt environments, this could enhance the hydrolysis of, for example, accumulated RNA, but the larger temperature differences also increase the movement of the interface by the recondensation of water, decreasing the time accumulated molecules spend dried at high temperatures. On the other side, the enhanced salt concentrations would trigger ribozymatic activity at the interface while the molecules in the bulk are protected by low salt concentrations from hydrolysis.

The system provided, in a single setting, a network of widely different reactions, connected by the fast diffusive transport of these molecules through water between different microbubbles. Ligation⁴¹ chemistry to drive replication could be made possible by

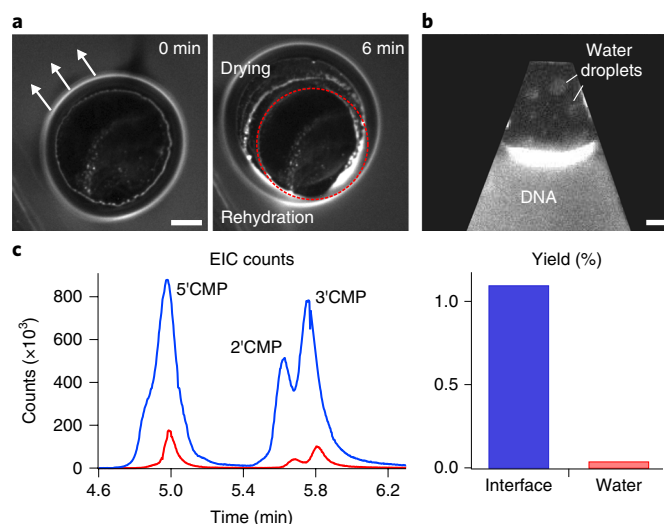


Fig. 7 | Dry-wet cycles and phosphorylation of nucleosides. **a**, Movement of bubbles in the chamber left behind dried material in the front that was rehydrated at the back side of the bubble, leading to dry-wet cycles (left, 0 min; right, 6 min). Dotted circle, the position of the accumulation region at $t = 0$ min. **b**, Using a hydrophobic cold side in the chamber increased the speed of dry-wet cycling by enhancing the formation of circular water droplets on the cold side. As they ‘rained’ into the bulk water, the interface moved periodically, triggering wet-dry cycles at the interface (Supplementary Video 6). **c**, Cytidine nucleosides were phosphorylated in the chamber shown in **b**. Over 12 h, the phosphorylation was, on average, 20 times more effective with the interfaces present (blue, $T_{\text{warm}} = 60$ °C, $T_{\text{cold}} = 37$ °C; Cytidine monophosphate (CMP) yield, $1.1 \pm 0.73\%$ from 12 repeats) compared to bulk water at 60 °C (red; CMP yield, $0.042 \pm 0.02\%$ from five repeats) or 50 °C ($0.056 \pm 0.0015\%$ from 15 repeats). Left, extracted-ion chromatogram (EIC) counts (EIC) for CMP peaks after HPLC-MS analysis. Right, yields measured by HPLC and ultraviolet detection. Scale bars, 100 μm (**a**), 250 μm (**b**).

the high local concentrations of substrates, while the drying process in close proximity would trigger the necessary phosphorylation^{37,38} and activation chemistry⁴² to drive the polymerization⁴³ of monomers. The dry-wet cycles were implemented continuously at the moving interface and molecules were retained at high concentrations at the interfaces after rehydration. In the adjacent water, molecules thermally cycled between warm and cold via laminar thermal convection.

At the interface, a strongly enhanced ribozymatic activity was demonstrated by the accumulation of the Hammerhead ribozyme. Its activity at the interface clearly dominates the turnover of substrate in the system, with up to 50% of the product being cleaved in a chamber with a gas interface (compared to 3.8% in the bulk). This can be attributed to higher local concentrations of ribozyme, substrate and MgCl_2 at the interface. Product strands are shuttled back into the bulk solution, where they are protected against the higher salt conditions. The mechanism thereby provides a way to enhance catalytic activity and increase the efficiency of RNA-catalysed processes.

Complex sequence phenotypes of RNA with several self-complementary sites have shown a sequence-selective formation of hydrogels at the interface. These hydrogels maintained a high local RNA concentration in water, an interesting setting to support efficient RNA catalysis^{44,45}, also because the high amount of hybridization in the hydrogels could protect oligonucleotides from hydrolysis⁴⁶, even in challenging salt concentrations.

In the presence of lipids such as oleic acid or DOPC, accumulation at the heated gas-water interface led to a continuous encapsulation of oligonucleotides into vesicle aggregates. The local DNA concentration inside these structures increased by a factor of up to 18 compared to the bulk solution. Folded RNA aptamers also accumulated inside the aggregates and dsDNA was shown to melt and re-anneal, demonstrating that aqueous phases readily exist inside the aggregates. As shown with DNase, the vesicles protected the encapsulated DNA from the bulk solution. The encapsulation of

oligonucleotides into lipid membranes is considered to be one of the key elements for more complex life and it has been suggested that lipids facilitated the assembly and polymerization of monomers^{47,48}. In the convection flow, the vesicles were subjected to temperature cycles and shear forces that led to vesicle fission. DOPC also produced vesicular structures, showing that modern phospholipids³⁹ could also accumulate and encapsulate oligonucleotides in the shown conditions.

For the prebiotic synthesis of RNA, crystallization at microbubbles would enable the purification of sugar-nucleobase precursors³⁴ and possibly also their chiral amplification by the enhanced growth of conglomerate RAO crystals³⁵. Interestingly, the sites of crystal formation later triggered, again, the formation of gas bubbles, showing a self-selection for crystallization conditions. If the seed crystal was homochiral, the subsequent bubble formation could accumulate and promote the assembly of more homochiral molecules at the same location. Finally, we found that the prebiotically important dry chemistry of nucleoside phosphorylation was enhanced by the gas interface: cytidine formed CMP 20 times more effectively compared to aqueous conditions.

To conclude, the experiments showed multiple modes of condensation, enrichment, accumulation and increased catalysis at heated gas microbubbles. This led to the physicochemical assembly and localization of prebiotic molecules—such as RNA precursors, lipids and ribozymes. We argue that this accumulation of molecules at a gas-water interface was a robust feature of natural microfluidic systems in porous volcanic rocks in aqueous environments, a setting likely to be ubiquitous on early Earth⁹. The simultaneous occurrence of six synergistic mechanisms for the accumulation and processing of prebiotic molecules, all operating in close proximity, fulfils the requirements for early life to connect a cascade of core reactions in the same non-equilibrium setting.

The setting presented here could therefore have largely helped in an informational polymer world, in which the first simple replicators were evolving. Following a synthesis of life’s first building

blocks, the accumulation at gas–water interfaces offers a mechanism to select polymers and enhance their catalytic activity. The shown continuous encapsulation dynamics of the accumulated molecules at the interface offers a pathway for the emergence of the cellular processes of life. Further experiments will test how this setting can host replication and selection towards early molecular evolution.

Methods

For the experiments in Fig. 2, a thermal chamber was sandwiched from a thin (>240 μm) layer of polyethylene terephthalate (PETG) plastic film deposited on a silicon wafer using a 3D printer (Ultimaker 2) in a funnel shape to facilitate a gas–water interface (Supplementary Fig. 1). A sapphire (Al_2O_3) block sealed the chamber. The chamber was filled through two 240- μm -thick borosilicate capillaries (Vitrocom). The system was annealed at 150 $^\circ\text{C}$ and the sapphire pressed down, fixing the thickness to 240 μm . Polydimethylsiloxane was deposited to seal the chamber. For most other experiments, chambers were built with an ultraviolet-curable resin (Photocentric 3D Daylight Resin, flexible, colour amber) through a master obtained by laser printing on a transparency film. Spacers between master and resin defined a chamber thickness of 150 μm or 250 μm . After illumination, sapphire windows were placed on top of the chambers and sealed with resin. Microfluidic access was provided through holes in the sapphire. Similarly, the chamber used for the phosphorylation and Hammerhead experiments was built by replacing the resin with a 254- μm -thick Teflon foil from which the structure was cut out using a cutting plotter. For the phosphorylation, additional Teflon foils were placed on the warm and cold sides of the chamber to mimic a hydrophobic surface. The temperature gradient was produced by heating the sapphire block through copper fixtures using heater cartridges for 3D printers and cooling the silicon side with a water bath. Temperature sensors and proportional–integral–derivative (PID) software maintained the temperatures. The temperature inside the chamber was calculated from the chamber geometry and known material constants with finite element methods (Comsol).

Fluorescence was measured with a fluorescence microscope (Zeiss Axio) through the transparent sapphire heating block using Mitutoyo infinity corrected long working distance objectives ($\times 2$ and $\times 10$) and a Zeiss Fluor $\times 5$ objective. The accumulation of DNA was detected with 200 nM FAM-labelled 132-base ssDNA in 0.1-fold PBS (see Supplementary Information for sequence). The silicon substrate was maintained at 10 $^\circ\text{C}$, with the copper heaters initially turned off. The experiment began with the copper heaters set to maintain a desired temperature. Background intensity levels were obtained from the non-fluorescing gas region. The bulk fluorescence signal was obtained from an area far from the free interface. The accumulated DNA fluorescence ratio was averaged perpendicular to the interface and divided by the bulk fluorescence. The flow was visualized with 200 nm FAM-labelled polystyrene beads in 0.1 \times PBS. The positions and velocities of the beads were tracked using ImageJ.

Simulation protocol. Simulations were performed using the finite element software COMSOL v4.4. The 2D model solved the convective heat equation, molecule diffusion equation and Navier–Stokes equations perpendicular to the contact line. Marangoni flows were established by implementing a stress boundary condition of the fluid velocity at the interface and by introducing a temperature-dependent surface tension. The water vapour concentration was simulated in the gas region above the interface by a diffusion–convection equation. The gas velocity was calculated from the temperature profile. Its velocity, combined with diffusion, caused an efficient net mass transport of vapour away from the interface into the gas bubble. By coupling a temperature-dependent vapour concentration boundary condition to the interface, a velocity boundary condition for water at the interface was imposed by the state equation of water. This resulted in capillary flow and evaporative mass transport of the water vapour, allowing vapour to enter and escape the water through evaporation or condensation. Thermophoresis was introduced via a thermophoretic drift term in the convection–diffusion equation describing the DNA concentration, and the Soret coefficients of the DNA were taken from experimental data⁴⁹. To incorporate the time lag of the heating process, the sapphire temperature was measured over time and incorporated as a polynomial function for T_{warm} .

The above simulation was solved over time, resulting in the time evolution of accumulation at the contact line. To fit the simulation to the observed data, the surface tension dependency on temperature was set within observed values⁵⁰ and fine-tuned as a free parameter using the bead tracking data. The geometry of the interface itself was adjusted to moderately tune the DNA accumulation dynamics. Concentrations were determined by averaging the top $\sim 10 \mu\text{m}$ of the meniscus at the hot side and comparing this to the simulated and experimental fluorescence, establishing a relation between fluorescence and average tip concentration.

To simulate the bubble shown in Supplementary Fig. 9, the simulation was transferred to an axial-symmetric geometry with spherical coordinates. The simulation around a small bubble was closed and gravity was pointed downwards. This removed the slow convection of water in the simulation, which, as expected from our modelling, did not change the accumulation characteristics. Ion and salt diffusion coefficients were taken from refs. 51–53.

Data availability

The data supporting the findings of this study are available within the paper and its Supplementary Information. Additional information and files are available from the corresponding author upon reasonable request. X-ray crystallographic data were also deposited at the Cambridge Crystallographic Data Centre (CCDC) under CCDC deposition no. 1847429.

Code availability

The complete details of both simulations are documented in the html report and mph simulation files in the Supplementary Information.

Received: 7 September 2018; Accepted: 21 June 2019;

Published online: 29 July 2019

References

- Schrödinger, E. *What is Life? The Physical Aspect of the Living Cell* (Cambridge Univ. Press, 1944).
- Cross, M. C. & Hohenburg, P. Pattern-formation outside of equilibrium. *Rev. Mod. Phys.* **65**, 851–1112 (1993).
- Bodenschatz, E., Pesch, W. & Ahlers, G. Recent developments in Rayleigh–Benard convection. *Annu. Rev. Fluid Mech.* **32**, 709–778 (2000).
- Fritts, D. C. & Alexander, M. J. Gravity wave dynamics and effects in the middle atmosphere. *Rev. Geophys.* **41**, 1003 (2003).
- Eaton, J. K. & Fessler, J. R. Preferential concentration of particles by turbulence. *Int. J. Multiph. Flow* **20**, 169–209 (1994).
- Götendorfer, A., Kruelle, C. A., Rehberg, I. & Svanšek, D. Localized subharmonic waves in a circularly vibrated granular bed. *Phys. Rev. Lett.* **97**, 198001 (2006).
- Chen, J. & Lopez, J. A. Interactions of platelets with subendothelium and endothelium. *Microcirculation* **12**, 235–246 (2005).
- Moore, W. B. & Webb, A. A. G. Heat-pipe Earth. *Nature* **501**, 501–505 (2013).
- Arndt, N. T. & Nisbet, E. G. Processes on the young Earth and the habitats of early life. *Annu. Rev. Earth Planet. Sci.* **40**, 521–549 (2012).
- Duhr, S. & Braun, D. Why molecules move along a temperature gradient. *Proc. Natl Acad. Sci. USA* **103**, 19678–19682 (2006).
- Baaske, P. et al. Extreme accumulation of nucleotides in simulated hydrothermal pore systems. *Proc. Natl Acad. Sci. USA* **104**, 9346–9351 (2007).
- Niether, D., Afanasenkau, D., Dhont, J. K. G. & Wiegand, S. Accumulation of formamide in hydrothermal pores to form prebiotic nucleobases. *Proc. Natl Acad. Sci. USA* **113**, 4272–4277 (2016).
- Kreysing, M., Keil, L., Lanzmich, S. & Braun, D. Heat flux across an open pore enables the continuous replication and selection of oligonucleotides towards increasing length. *Nat. Chem.* **7**, 203–208 (2015).
- Mast, C. B., Schink, S., Gerland, U. & Braun, D. Escalation of polymerization in a thermal gradient. *Proc. Natl Acad. Sci. USA* **110**, 8030–8035 (2013).
- Morasch, M., Braun, D. & Mast, C. B. Heat-flow-driven oligonucleotide gelation separates single-base differences. *Angew. Chem. Int. Ed.* **55**, 6676–6679 (2016).
- Keil, L. M. R., Möller, F. M., Kieß, M., Kudella, P. W. & Mast, C. B. Proton gradients and pH oscillations emerge from heat flow at the microscale. *Nat. Commun.* **8**, 1897 (2017).
- Budin, I., Bruckner, R. J. & Szostak, J. W. Formation of protocell-like vesicles in a thermal diffusion column. *J. Am. Chem. Soc.* **131**, 9628–9629 (2009).
- Lerman, L. Potential role of bubbles and droplets in primordial and planetary chemistry. *Orig. Life Evol. Biosph.* **16**, 201–202 (1986).
- Ariga, K. & Hill, J. P. Monolayers at air–water interfaces: from origins-of-life to nanotechnology. *Chem. Rec.* **11**, 199–211 (2011).
- Eickbush, T. H. & Moudrianakis, E. N. A mechanism for the entrapment of DNA at an air–water interface. *Biophys. J.* **18**, 275–288 (1977).
- Griffith, E. C. & Vaida, V. In situ observation of peptide bond formation at the water–air interface. *Proc. Natl Acad. Sci. USA* **109**, 15697–15701 (2012).
- Deegan, R. D. et al. Capillary flow as the cause of ring stains from dried liquid drops. *Nature* **389**, 827–829 (1997).
- Smith, K. A. On convective instability induced by surface-tension gradients. *J. Fluid Mech.* **24**, 401–414 (1966).
- Batchelor, G. K. *An Introduction to Fluid Dynamics* (Cambridge Univ. Press, 1973).
- Deegan, R. D. Pattern formation in drying drops. *Phys. Rev. E* **61**, 475–485 (2000).
- Larson, R. G. Transport and deposition patterns in drying sessile droplets. *AIChE J.* **60**, 1538–1571 (2014).
- Savino, R., Paterna, D. & Favaloro, N. Buoyancy and marangoni effects in an evaporating drop. *J. Thermophys. Heat Transf.* **16**, 562–574 (2002).
- Drobot, B. et al. Compartmentalised RNA catalysis in membrane-free coacervate protocells. *Nat. Commun.* **9**, 3643 (2018).
- Weinberg, M. S. & Rossi, J. J. Comparative single-turnover kinetic analyses of trans-cleaving hammerhead ribozymes with naturally derived non-conserved sequence motifs. *FEBS Lett.* **579**, 1619–1624 (2005).

30. Dahm, S. C. & Uhlenbeck, O. C. Role of divalent metal ions in the hammerhead RNA cleavage reaction. *Biochemistry* **30**, 9464–9469 (1991).
31. Zhu, T. F. & Szostak, J. W. Coupled growth and division of model protocell membranes. *J. Am. Chem. Soc.* **131**, 5705–5713 (2009).
32. Budin, I. & Szostak, J. W. Physical effects underlying the transition from primitive to modern cell membranes. *Proc. Natl Acad. Sci. USA* **108**, 5249–5254 (2011).
33. Filonov, G. S., Moon, J. D., Svensen, N. & Jaffrey, S. R. Broccoli: rapid selection of an RNA mimic of green fluorescent protein by fluorescence-based selection and directed evolution. *J. Am. Chem. Soc.* **136**, 16299–16308 (2014).
34. Islam, S., Bučar, D.-K. & Powner, M. W. Prebiotic selection and assembly of proteinogenic amino acids and natural nucleotides from complex mixtures. *Nat. Chem.* **9**, 584–589 (2017).
35. Anastasi, C., Crowe, M., Powner, M. W. & Sutherland, J. D. Direct assembly of nucleoside precursors from two- and three-carbon units. *Angew. Chem. Int. Ed.* **45**, 6176–6179 (2006).
36. Jones, S. F., Evans, G. M. & Galvin, K. P. Bubble nucleation from gas cavities—a review. *Adv. Colloid Interface Sci.* **80**, 27–50 (1999).
37. Powner, M. W., Gerland, B. & Sutherland, J. D. Synthesis of activated pyrimidine ribonucleotides in prebiotically plausible conditions. *Nature* **459**, 239–242 (2009).
38. Lohrmann, R. & Orgel, L. E. Urea–inorganic phosphate mixtures as prebiotic phosphorylating agents. *Science* **171**, 490–494 (1971).
39. Gibard, C., Bhowmik, S., Karki, M., Kim, E.-K. & Krishnamurthy, R. Phosphorylation, oligomerization and self-assembly in water under potential prebiotic conditions. *Nat. Chem.* **10**, 2012–2017 (2017).
40. Chow, Y. T. F., Maitland, G. C. & Trusler, J. P. M. Interfacial tensions of the (CO₂+N₂+H₂O) system at temperatures of (298 to 448) K and pressures up to 40 MPa. *J. Chem. Thermodyn.* **93**, 392–403 (2016).
41. Sosson, M. & Richter, C. Enzyme-free genetic copying of DNA and RNA sequences. *Beilstein J. Org. Chem.* **14**, 603–617 (2018).
42. Forsythe, J. G. et al. Ester-mediated amide bond formation driven by wet–dry cycles: a possible path to polypeptides on the prebiotic earth. *Angew. Chem. Int. Ed.* **54**, 9871–9875 (2015).
43. Morasch, M., Mast, C. B., Langer, J. K., Schilcher, P. & Braun, D. Dry polymerization of 3',5'-cyclic GMP to long strands of RNA. *ChemBioChem* **15**, 879–883 (2014).
44. Vaidya, N. et al. Spontaneous network formation among cooperative RNA replicators. *Nature* **491**, 72–77 (2012).
45. Mutschler, H., Wochner, A. & Holliger, P. Freeze–thaw cycles as drivers of complex ribozyme assembly. *Nat. Chem.* **7**, 502–508 (2015).
46. Soukup, G. A. & Breaker, R. R. Relationship between internucleotide linkage geometry and the stability of RNA. *RNA* **5**, 1308–1325 (1999).
47. Topozini, L., Dies, H., Deamer, D. W. & Rheinstädter, M. C. Adenosine monophosphate forms ordered arrays in multilamellar lipid matrices: insights into assembly of nucleic acid for primitive life. *PLoS One* **8**, e62810 (2013).
48. Rajmani, S. et al. Lipid-assisted synthesis of RNA-like polymers from mononucleotides. *Orig. Life Evol. Biosph.* **38**, 57–74 (2008).
49. Reineck, P., Wienken, C. J. & Braun, D. Thermophoresis of single stranded DNA. *Electrophoresis* **31**, 279–286 (2010).
50. Vargaftik, N. B., Volkov, B. N. & Voljak, L. D. International tables of the surface tension of water. *J. Phys. Chem. Ref. Data* **12**, 817–820 (1983).
51. Lide, D. R. *CRC Handbook of Chemistry and Physics* (CRC Press, 2001).
52. Li, Y. & Gregory, S. Diffusion of ions in sea water and in deep-sea sediments. *Geochim. Cosmochim. Acta* **33**, 703–714 (1974).
53. Fell, C. J. D. & Hutchison, H. P. Diffusion coefficients for sodium and potassium chlorides in water at elevated temperatures. *J. Chem. Eng. Data* **16**, 427–429 (1971).

Acknowledgements

The authors thank L. Keil for help with data analysis. Financial support from the Simons Foundation (318881 to M.W.P. and 327125 to D.B.), the German Research Foundation (DFG) through CRC/SFB 235 Project P07 and SFB 1032 Project A04, DFG Grant BR2152/3-1 and the US–German Fulbright Program is acknowledged. H.M. is supported by the MaxSynBio consortium, which is jointly funded by the Federal Ministry of Education and Research of Germany and the Max Planck Society. H.M. and K.L.V. are supported by the Volkswagen Initiative 'Life'—A Fresh Scientific Approach to the Basic Principles of Life. A.K. is supported by a DFG fellowship through the Graduate School of Quantitative Biosciences Munich.

Author contributions

M.M., J.L., C.F.D., A.K., A.I. and Ph.S. performed the experiments. M.M., J.L., K.L.V., S.I., B.S., D.B.D., H.M., Pe.S., M.W.P., C.B.M. and D.B. conceived and designed the experiments. M.M., J.L., K.L.V., S.I., M.K.C., H.M., M.W.P. and D.B. analysed the data. M.M., J.L. and D.B. wrote the paper. All authors discussed the results and commented on the manuscript.

Competing interests

The authors declare no competing interests.

Additional information

Supplementary information is available for this paper at <https://doi.org/10.1038/s41557-019-0299-5>.

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